

Global surface of section and Hamiltonian dynamics on 2-disk

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The goal of this talk is to prove the following result.

Theorem (Abbondandolo-Bramham-Hryniewicz-Salomão, 2018)

There exists a contact form α on S^3 such that

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(1) The conclusion above does *not* hold for $\alpha = \alpha_0$, the standard contact structure of S^3 . In fact, $T_{\min} = \pi$ and $\text{vol}_{\alpha \wedge d\alpha}(S^3) = \pi^2$ (Exercise).

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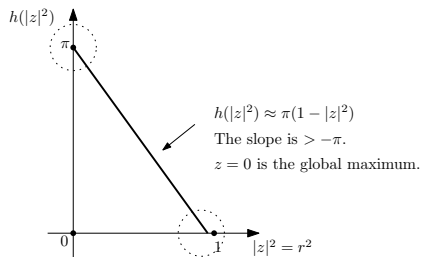
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$$\text{vol}_{\alpha \wedge d\alpha}(S^3) = \pi^2 + \text{Cal}(\phi). \quad (1)$$

Moreover, this ϕ is special since $\text{Cal}(\phi) < 0$ but $T_{\min}(\alpha) \geq \pi$.

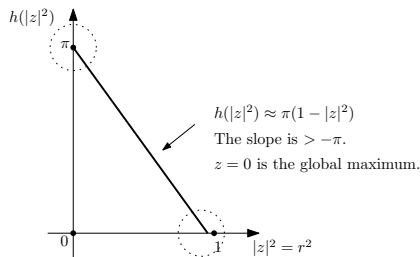
Radial symmetric Hamiltonian on \mathbb{D}

- Consider the autonomous Hamiltonian $H : \mathbb{D} \rightarrow \mathbb{R}$ such that $H(z) = h(|z|^2)$ where $h : [0, 1] \rightarrow \mathbb{R}$ is defined in the following picture (where circle regions are smoothed),



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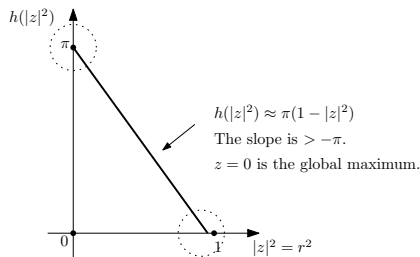
$$(2) \text{Cal}(\phi_H^1) \approx \pi^2.$$

In general,

$$\text{Cal}(\phi_H^1) = 4\pi \int_0^1 rh(r^2)dr.$$

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- Observe that fixed points lie in neighborhood of $\partial\mathbb{D}$ and $z = 0$.

Action function

Given a radial symmetric Hamiltonian function $H : \mathbb{D} \rightarrow \mathbb{R}$, where $H(z) = h(|z|^2)$. Consider **action function** $\sigma : \mathbb{D} \rightarrow \mathbb{R}$ defined by

$$\sigma(z) = h(|z|^2) - |z|^2 h'(|z|^2).$$

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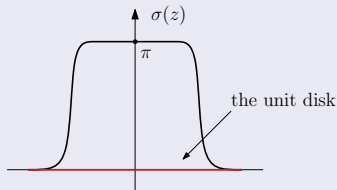
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Example

For H chosen earlier, action function is the following graph.



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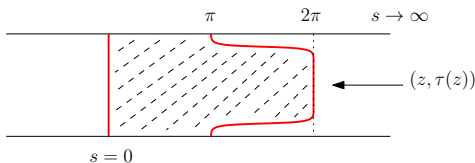
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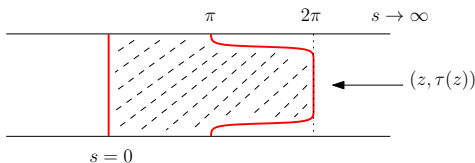


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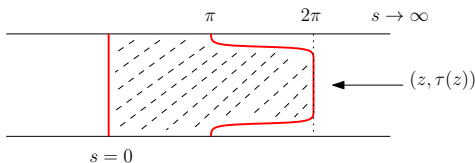
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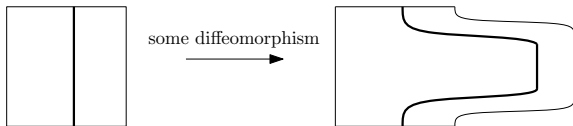
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- Observe that $t^*(\lambda + ds) = \lambda + ds$ (Exercise). Therefore, M admits a well-defined contact form denoted by η .
- Near the boundary, $M = U \times \mathbb{R} / \pi\mathbb{Z}$ where U is a neighborhood of $\partial\mathbb{D}$ and $\eta = \lambda + ds$.

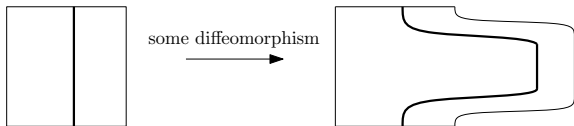
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The desired contact form β is on the solid torus $\mathbb{D} \times \mathbb{R}/\pi\mathbb{Z}$. There exists some diffeomorphism $f : \mathbb{D} \times \mathbb{R}/\pi\mathbb{Z} \rightarrow M$ as in the following picture.



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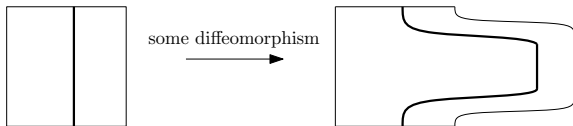
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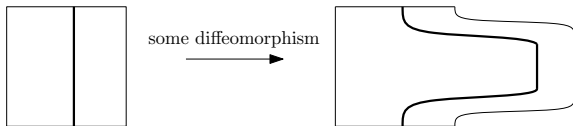
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Example

One can compute that $\text{vol}_{\beta \wedge d\beta}(\mathbb{D} \times \mathbb{R}/\pi\mathbb{Z}) = \text{vol}_{\eta \wedge d\eta} M \simeq 2\pi^2$. Then, up to small gaps, we get

$$\text{vol}_{\beta \wedge d\beta}(\mathbb{D} \times \mathbb{R}/\pi\mathbb{Z}) = \pi^2 + \text{Cal}(\phi_H^1).$$

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$$g(r, \theta, s) = \left(re^{i(\theta+2s)}, \sqrt{1-r^2}e^{2is} \right)$$

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- $\text{vol}_{\alpha \wedge d\alpha}(S^3) = \text{vol}_{\beta \wedge d\beta}(\mathbb{D} \times \mathbb{R}/\pi\mathbb{Z})$, which implies the volume-Calabi equation (1). Moreover, $T_{\min}(\alpha) = \pi$ (Exercise).

A cheap variation

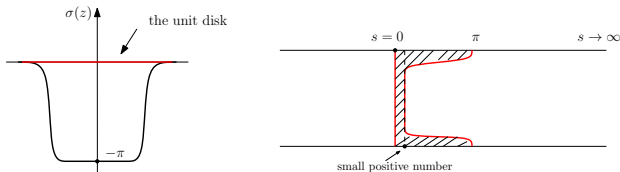
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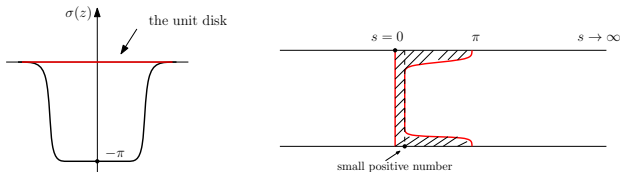
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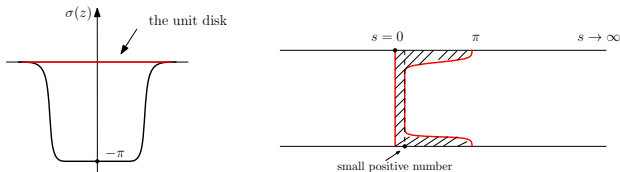
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The equation (1) holds, but unfortunately $T_{\min}(\alpha)$ is very small!

- Need a new example to satisfy both $\text{Cal}(\phi) < 0$ and $T_{\min}(\alpha) \geq \pi$.

Sinkhole (cf. Usher's Banach-Mazur distance paper)

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- Trick: Fix an ω -preserving diffeomorphism ψ on $\mathring{\mathbb{D}}$ such that $\psi(z) = z + z_0$ on $\mathbb{D}_{(0, r_0)}$, and consider $G = H \circ \psi$.

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• Trick: Fix an ω -preserving diffeomorphism ψ on $\mathring{\mathbb{D}}$ such that $\psi(z) = z + z_0$ on $\mathbb{D}_{(0, r_0)}$, and consider $G = H \circ \psi$. Then the morphism $\phi_G^1 = \psi^{-1} \circ \phi_H^1 \circ \psi$ rotates the disk $\mathbb{D}_{(0, r_0)}$.

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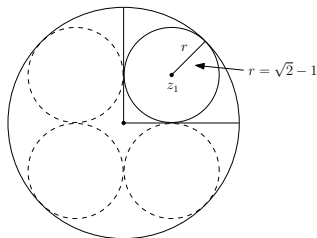
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An estimation: $|u \circ (\psi^{-1} \circ \phi) - u \circ \psi^{-1}| \leq r_0 |z_0| \leq r_0$, which implies that

$$-k - r_0 \leq \sigma_\phi \leq r_0. \quad (2)$$

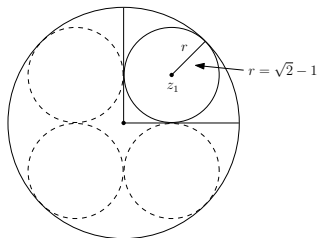
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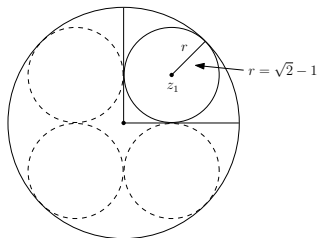


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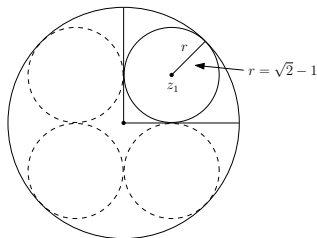
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Denote

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Then we can compute

$$\begin{aligned}\sigma_\phi &= \sigma_{\phi^+} \circ \left(\prod_{i=1}^4 \phi_i^- \right) + \sigma_{\prod_{i=1}^4 \phi_i^-} \\ &= \sigma_{\phi^+} \circ \left(\prod_{i=1}^4 \phi_i^- \right) + \sum_{i=1}^4 \sigma_{\phi_i^-}.\end{aligned}$$

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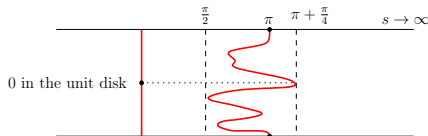
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A schematic picture of Bramham's solid torus (denoted earlier by M) is the following.



Final step: estimation of $T_{\min}(\alpha)$

Closed Reeb orbit comes from two cases: (1) fixed point of ϕ on \mathbb{D} ; (2) k -periodic points of ϕ on \mathbb{D} .

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The upshot is:

$$\left\{ \begin{array}{l} \text{vol}_{\alpha \wedge d\alpha}(S^3) = \pi^2 + \text{Cal}(\phi) \\ \text{Cal}(\phi) < 0 \\ T_{\min}(\alpha) \geq \pi \end{array} \right. \Rightarrow T_{\min}(\alpha)^2 > \text{vol}_{\alpha \wedge d\alpha}(S^3).$$

Outlook

- One can re-do the computation above by starting with a very narrow sector of \mathbb{D} and rotate it making it symmetric. Then the resulting ϕ will be C^0 -close to $1_{\mathbb{D}}$ (and α is C^0 -closed to α_0).

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Its proof runs our construction backwards, i.e., constructing an ω -preserving diffeomorphism on \mathbb{D} from a contact form on S^3 .

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- Usher's tube construction: given a Liouville domain (W, λ) and an autonomous Hamiltonian function $H : W \rightarrow \mathbb{R}$ satisfying certain "positivity condition", the following construction results in another Liouville domain,

$$W_H := \{(w, z) \in W \times \mathbb{C} \mid \pi|z|^2 \leq H(w)\}.$$